

MSDL presentation

A Brief Introduction to Statistical Mechanics

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Overview

- The Hamiltonian, the Hamiltonian equations of motion; simple harmonic oscillator.
- Thermodynamics; internal energy and free energy.
- Kinetic theory; phase space and distribution functions.
- Statistical mechanics: ensembles, canonical partition function; the ideal gas.
- Concluding remarks.

Equations of Motion

- Newton's second law: $\mathbf{F} = m \mathbf{a}$.
- Equations of motion: explicit expressions for \mathbf{F} , \mathbf{a} .
- Different formulations give a recipe to arrive at this: **Lagrangian** and **Hamiltonian** approaches.
- In Classical Mechanics, the Lagrangian formulation is most common – leading to Lagrange's equations of motion.
- The Hamiltonian formulation leads to Hamilton's equations of motion.
- In Statistical Mechanics and Quantum Mechanics it is more convenient to use the Hamiltonian formalism.

The Hamiltonian

- The Hamiltonian \mathcal{H} is defined as follows:

$$\mathcal{H} = T + V.$$

T : kinetic energy of the system,

V : potential energy of the system.

- For most systems, \mathcal{H} is just the total energy E of the system.
- Knowing \mathcal{H} , we can write down the equations of motion.
- In Classical Physics, $\mathcal{H} \equiv E$, a scalar quantity.
- In Quantum Mechanics, \mathcal{H} is an operator.

Generalized Coordinates

- The concept of a 'coordinate' is extended, appropriate variables can be used. For each generalized coordinate, there is a corresponding generalized momentum (canonically conjugate). Notation:

q_i : generalized coordinates;

p_i : generalized momenta.

- For example,
 - $\{q_i\} \equiv$ Cartesian coordinates $\{x, y, z\}$ for a free particle.
 - $\{q_i\} \equiv$ the angle θ for a simple pendulum.
- Can transform between the regular coordinates and the generalized coordinates.
- Useful in dealing with constraints.

Generalized Coordinates

- Hamiltonian $\mathcal{H} = \mathcal{H}(\{q_i\}, \{p_i\}, t) = \mathcal{H}(q, p, t)$.
- The generalized coordinates and momenta need not correspond to the usual spatial coordinates and momenta.
- In general for N particles, we have $3N$ coordinates and $3N$ momenta.
- Coordinates and momenta are considered 'conjugate' quantities - they are on an equal footing in the Hamiltonian formalism.
- Coordinates and momenta together define a **phase space** for the system.

Equations of Motion

In the Hamiltonian formalism, the equations of motion are given by:

$$\frac{\partial \mathcal{H}}{\partial q_i} = -\dot{p}_i;$$

$$\frac{\partial \mathcal{H}}{\partial p_i} = \dot{q}_i;$$

Newton's second law!

Equations of Motion

- If there is no explicit time dependence in the Hamiltonian, then

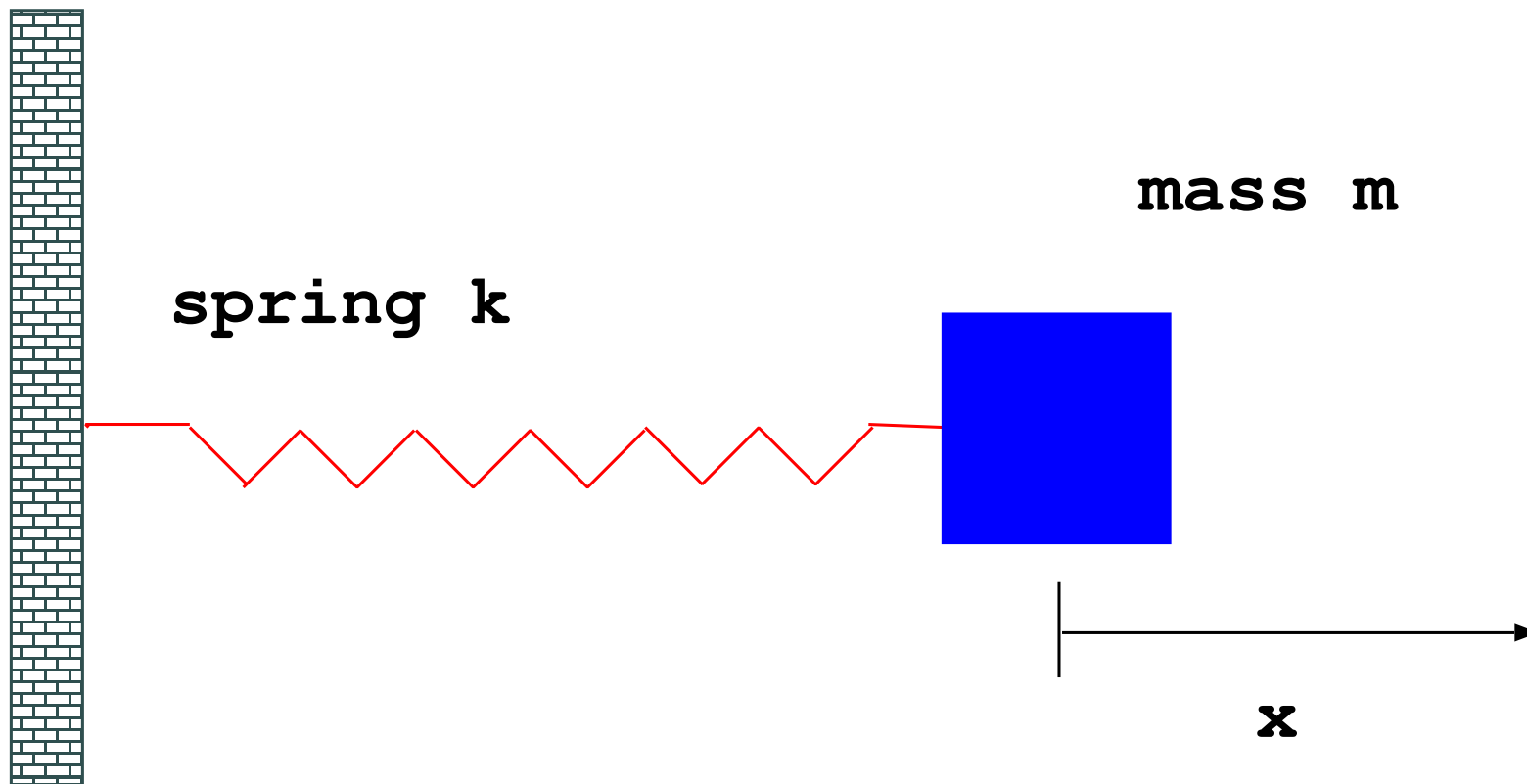
$$\mathcal{H}(q, p, t) = \mathcal{H}(q, p);$$

$$\frac{d\mathcal{H}}{dt} = \frac{\partial\mathcal{H}}{\partial t} = 0.$$

\implies The energy of the system does not change with time, and it is conserved.

- Symmetries of the Hamiltonian \implies conserved quantities.

Simple Harmonic Oscillator (1-D)



Example: Simple Harmonic Oscillator (1-D)

- Consider a block of mass m connected by a spring with spring constant k .
- Its displacement is given by x , has velocity v , momentum $p = mv$.
- One generalized co-ordinate x , its conjugate momentum p .
- The total energy : kinetic energy + potential energy

$$E = \frac{1}{2}mv^2 + \frac{1}{2}kx^2.$$

Simple Harmonic Oscillator

The Hamiltonian:

$$E = \frac{1}{2} m v^2 + \frac{1}{2} k x^2.$$
$$\mathcal{H} = \frac{p^2}{2m} + \frac{1}{2} k x^2.$$

Hamilton's equations of motion:

$$\frac{\partial \mathcal{H}}{\partial x} = -\dot{p};$$

$$\frac{\partial \mathcal{H}}{\partial p} = \dot{x}.$$

Simple Harmonic Oscillator

$$\begin{aligned}kx &= -\dot{p}; \\ \frac{p}{m} &= \dot{x};\end{aligned}$$

Rearrange, two first order ODEs:

$$\begin{aligned}\frac{dp}{dt} &= -kx; \quad \text{Newton's Law} \\ \frac{dx}{dt} &= \frac{p}{m}.\end{aligned}$$

Can get the usual second order ODE:

$$\frac{d^2x}{dt^2} + \frac{k}{m}x = 0.$$

Thermodynamics

- Thermodynamics: phenomenological and empirical.
- A thermodynamic system: any macroscopic system.
- Thermodynamic parameters (state variables): measurable quantities such as pressure P , volume V , temperature T , magnetic field H .
- A thermodynamic state is specified by particular values of $P, V, T, H \dots$
- An **equation of state**: a functional relation between the state variables. Example: for an ideal gas, $PV = nRT$.
- Other thermodynamic quantities: internal energy E , entropy S , specific heats C_V, C_P (response functions).

Thermodynamics

- At equilibrium, observe **average** behaviour.
- The internal energy E = average total energy of the system : $\langle \mathcal{H} \rangle$.
- **Intensive** and **Extensive** variables.
- Intensive variables do not depend on the size of the system.
Examples: pressure P , temperature T , chemical potential μ .
- Extensive variables depend on the size of the system. Examples: the total number of particles N , volume V , internal energy E , entropy S .
- The first law of thermodynamics: energy conservation:
change in internal energy = heat supplied - work done by the system.
Mathematically:

$$dE = T dS - P dV + \mu dN .$$

Thermodynamics

- T, P, μ are **generalized forces**, intensive. Each associated with an extensive variable, such that
change in internal energy = generalized force \times change in extensive variable.
- $\{T, S\}, \{-P, V\}, \{\mu, N\}$ are **conjugate variable pairs**.
- Thermodynamic potentials: analogous to the mechanical potential energy. Energy available to do work – ‘free energy’. The free energy is minimized, depending on the conditions.

Thermodynamics

- Helmholtz free energy:

$$F = E - TS;$$

$$dF = dE - T dS - S dT$$

$$= T dS - P dV + \mu dN - T dS - S dT$$

$$dF = -P dV - S dT + \mu dN.$$

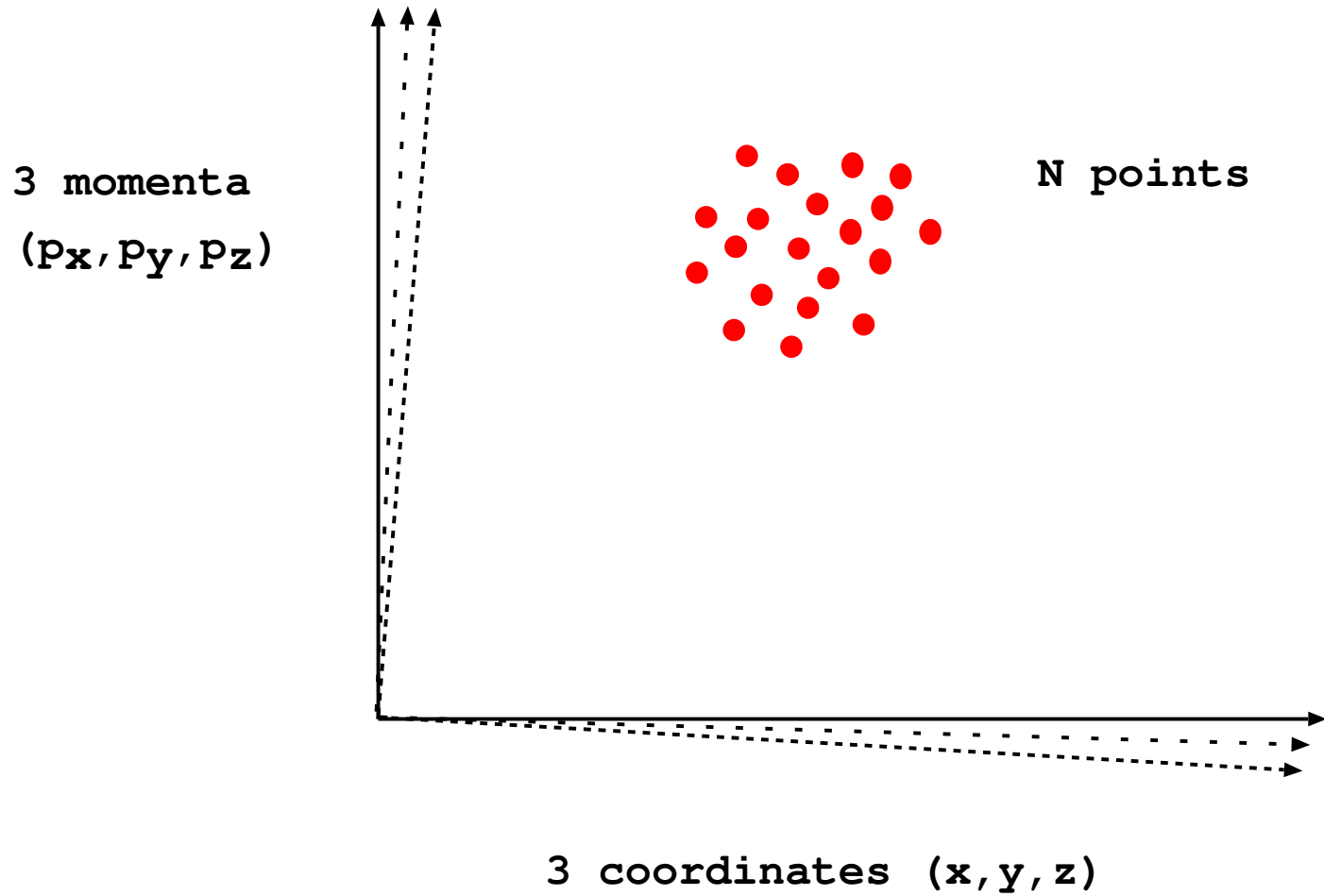
- Like E , the potentials contain all thermodynamic information.
- Equation of state from the free energy – relates state variables:

$$P(V, T, N) = - \left. \frac{\partial F}{\partial V} \right|_{T, N}.$$

Kinetic Theory

- A dilute gas of a large number N of molecules in a volume V .
- The temperature T is high, the density is low.
- The molecules interact via collisions.
- An isolated system will always reach equilibrium by minimizing its energy. At equilibrium its energy does not change with time.
- Consider, for each molecule, $\{\mathbf{r}, \mathbf{p}\}$: 3 spatial coordinates, 3 momenta.
- Each particle corresponds to a point in a 6-D phase space. The system as a whole can be represented as N points.
- Not interested in the detailed behaviour of each molecule.

Kinetic Theory – 6-D Phase Space



Kinetic Theory

- Define a **distribution function** $f(\mathbf{r}, \mathbf{p}, t)$ so that

$$f(\mathbf{r}, \mathbf{p}, t) d\mathbf{r} d\mathbf{p}$$

gives the number of molecules at time t lying within a volume element $d\mathbf{r}$ about \mathbf{r} and with momenta in a momentum-space volume $d\mathbf{p}$ about \mathbf{p} .

- $f(\mathbf{r}, \mathbf{p}, t)$ is just the **density of points** in phase space.
- Assume that f is a smoothly varying continuous function, so that:

$$\int f(\mathbf{r}, \mathbf{p}, t) d\mathbf{r} d\mathbf{p} = N.$$

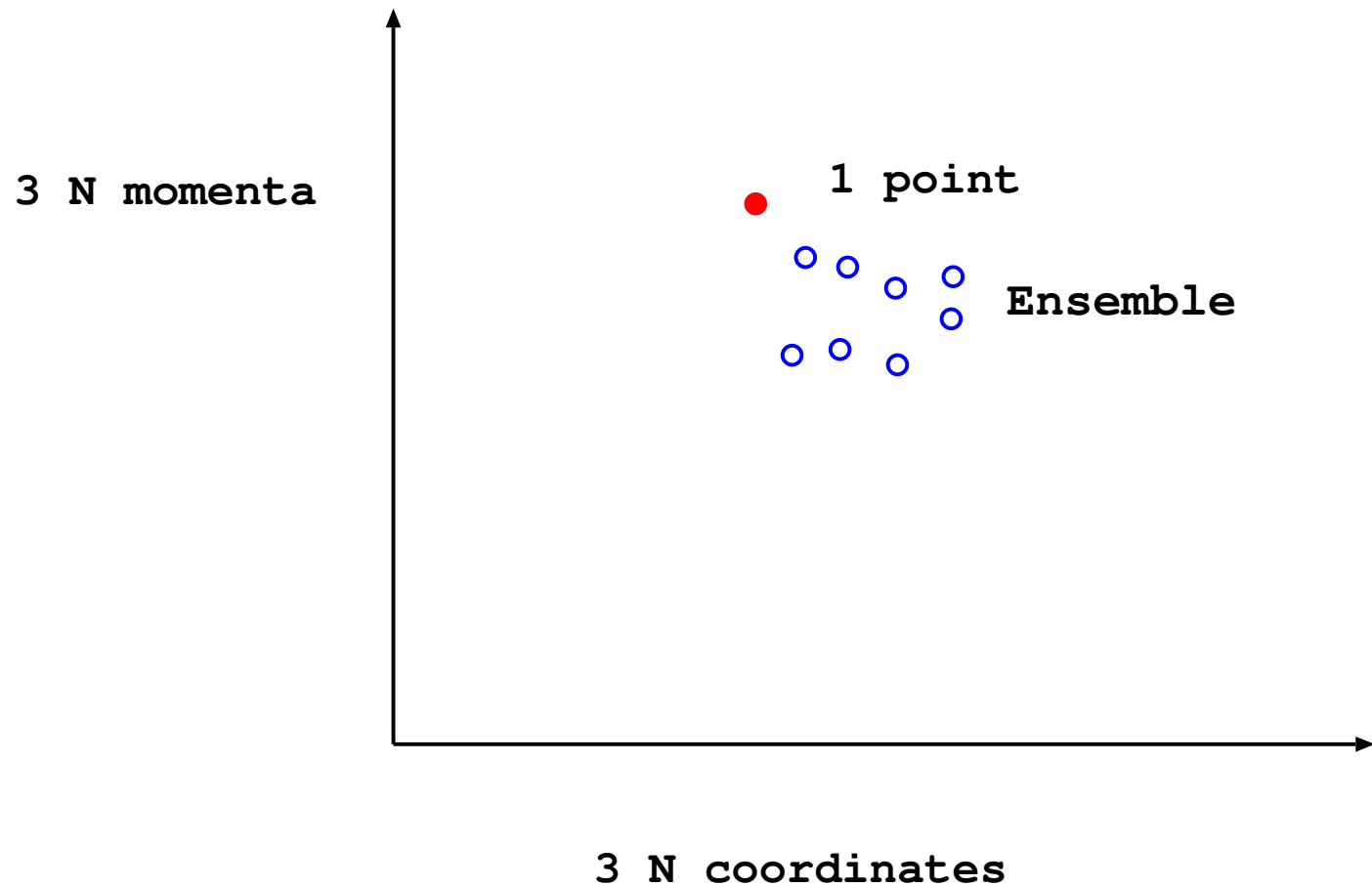
Kinetic Theory

- Problem of Kinetic Theory: Find f for given kinds of collisions (binary, for example) – can be considered to be a form of interaction.
- Explicit expressions for pressure, temperature; temperature is a measure of the average kinetic energy of the molecules.
- The limiting form of f as $t \rightarrow \infty$ will yield all the equilibrium properties for the system, and hence the thermodynamics.
- Maxwell–Boltzmann distribution of speeds for an ideal gas at equilibrium at a given temperature: Gaussian.
- Find the time-evolution equation for f - the equation of motion in phase space.
- Very messy!

Ensembles and Statistical Mechanics

- Gibbs introduced the concept of an **ensemble**.
- Earlier: N particles in the 6-D phase space.
- Ensemble: A **single** point in $6N$ -dimensional phase space represents a given configuration.
- A given macrostate with $\{E, T, P, V, \dots\}$ corresponds to an ensemble of microstates: ‘snapshots’ of the system at different times.
- Ergodic hypothesis: Given enough time, the system explores all possible points in phase space.

Ensemble Theory – $6N$ -D Phase Space



Ensembles and Statistical Mechanics

- Central idea: replace **time averages** by **ensemble averages**.
- Time Average: For any quantity $\phi(\mathbf{r}, \mathbf{p}, t)$:

$$\langle \phi(\mathbf{r}, \mathbf{p}, t) \rangle = \frac{1}{\tau} \int_0^\tau \phi(\mathbf{r}, \mathbf{p}, t) dt.$$

- Ensemble Average: For $\tau \rightarrow \infty$:

$$\langle \phi(\mathbf{r}, \mathbf{p}, t) \rangle = \int \rho(\mathbf{r}, \mathbf{p}) \phi(\mathbf{r}, \mathbf{p}, t) d\mathbf{r} d\mathbf{p}.$$

$$d\mathbf{r} = d\mathbf{r}_1 d\mathbf{r}_2 \dots d\mathbf{r}_N;$$

$$d\mathbf{p} = d\mathbf{p}_1 d\mathbf{p}_2 \dots d\mathbf{p}_N.$$

Ensembles

- $\rho(\mathbf{r}, \mathbf{p})$: probability density,
 $\rho(\mathbf{r}, \mathbf{p}) d\mathbf{r} d\mathbf{p}$: probability of finding the system in a volume element $[d\mathbf{r} d\mathbf{p}]$ around (\mathbf{r}, \mathbf{p}) .
- Different ensembles: **microcanonical**, **canonical** and **grand canonical** ensembles.
- Microcanonical ensemble: isolated systems with fixed energy and number of particles, no exchange of energy or particles with the outside world. Not very useful !

Ensembles

- Canonical ensemble: Energy is not fixed, can exchange E with a reservoir; N fixed.
- Grand canonical ensemble: Both energy and N can vary.
- In the canonical ensemble, the probability of a given configuration with energy E (corresponding to Hamiltonian \mathcal{H}) :

$$p_c = \frac{e^{-\beta \mathcal{H}(\mathbf{r}, \mathbf{p})}}{Z_N(V, T)} .$$

- $\beta = 1/k_B T$, k_B : Boltzmann constant = R/N_A .
- $e^{-\beta \mathcal{H}(\mathbf{r}, \mathbf{p})}$: Boltzmann factor or Boltzmann weight.

The Canonical Partition Function

$$Z_N(V, T) = \frac{1}{N! h^{3N}} \int d\mathbf{r} d\mathbf{p} e^{-\beta \mathcal{H}(\mathbf{r}, \mathbf{p})}$$

h : Planck's constant.

- Z_N : phase space volume, each volume element weighted with the Boltzmann factor.
- From Z_N , we can calculate various thermodynamic quantities.

The Canonical Partition Function

- For example: internal energy E is given by the ensemble average $\langle \mathcal{H} \rangle$ of the Hamiltonian:

$$\begin{aligned} E &= \langle \mathcal{H} \rangle \\ &= \frac{1}{Z_N(V, T)} \int d\mathbf{r} d\mathbf{p} \mathcal{H} e^{-\beta \mathcal{H}(\mathbf{r}, \mathbf{p})} \end{aligned}$$

- Also, more conveniently,

$$E = - \left. \frac{\partial \ln Z_N(V, T)}{\partial \beta} \right|_{N, V} .$$

The Canonical Partition Function

Can show that it is related to the Helmholtz free energy F :

$$F = E - T S$$

$$Z_N(V, T) = e^{-\beta F}$$

$$F = -k_B T \ln Z_N(V, T).$$

Example: The Ideal Gas

- One of the simplest statistical systems: a gas of N non-interacting particles, each of mass m , in a volume V at temperature T .
- The Hamiltonian \mathcal{H} = Kinetic Energy:

$$\mathcal{H} = \sum_{i=1}^N \frac{\mathbf{p}_i^2}{2m} = \sum_{i=1}^N \frac{p_i^2}{2m}. \quad (1)$$

- The canonical partition function:

$$Z_N(V, T) = \frac{1}{N! h^{3N}} \int d\mathbf{r} d\mathbf{p} e^{-\beta \mathcal{H}(\mathbf{r}, \mathbf{p})}$$
$$Z_N(V, T) = \frac{1}{N! h^{3N}} \int d\mathbf{r} d\mathbf{p} \exp \left(- \sum_{i=1}^N \frac{p_i^2}{2m} \right).$$

Example: The Ideal Gas

$$Z_N = \frac{V^N}{N! h^{3N}} \left(\int_{-\infty}^{\infty} dp e^{-\beta p^2 / 2m} \right)^{3N}$$

$$Z_N = \frac{V^N}{N! h^{3N}} \left(\sqrt{2\pi m / \beta} \right)^{3N}$$

$$= \frac{V^N}{N! h^{3N}} (2\pi m k_B T)^{3N/2}$$

$$\ln Z_N = N \ln V + \frac{3}{2} N [\ln(2\pi m) - \ln \beta] - \ln N! - 3N \ln h.$$

The Ideal Gas

- Can calculate macroscopic thermodynamical quantities.
- Average (kinetic) energy of the gas:

$$\begin{aligned} E &= - \frac{\partial \ln Z_N}{\partial \beta} \\ &= - \frac{\partial(-\frac{3}{2} N \ln \beta)}{\partial \beta} \\ &= \frac{3}{2} \frac{N}{\beta} \\ E &= \frac{3}{2} N k_B T . \end{aligned}$$

- Average (kinetic) energy per particle: $\frac{E}{N} = \frac{3}{2} k_B T$.

The Ideal Gas

$$\ln Z_N = N \ln V + \frac{3}{2} N [\ln(2\pi m) - \ln \beta] - \ln N! - 3N \ln h.$$

The Helmholtz free energy:

$$F = -k_B T \ln Z_N.$$

The ideal gas equation of state:

$$P(V, T, N) = - \left. \frac{\partial F}{\partial V} \right|_{T, N}.$$

The Ideal Gas

Equation of state:

$$\begin{aligned} P &= - \frac{\partial(-k_B T N \ln V)}{\partial V} \\ &= \frac{Nk_B T}{V} \end{aligned}$$

$$PV = nN_A k_B T$$

$$PV = nRT$$

The Partition Function

- The partition function approach is very important because of its success.
- A whole range of equilibrium phenomena can be understood this way.
- Recipe: Write down the Hamiltonian - mostly interested in the potential energy term:

$$V = V_{ext} + V_{int}$$

- The V_{int} term includes all the interactions: two-body, three-body ...
- Get the free energy from the partition function.
- Minimize the free energy: the equilibrium 'ground state' of the system.

Extensions ?

- The Hamiltonian can be either discrete or continuous.
- Continuous systems: concept of **fields** (example: density field, magnetization field, field of interactions).
- When dealing with fields, the Hamiltonian, and the free energy, become **functionals** :

$$\mathcal{H} = \mathcal{H} [\phi(\mathbf{r}, t)]$$

- Use field theory techniques and variational calculus.
- Non-equilibrium: time-dependent Hamiltonian, dissipation \implies transport properties; failure of equilibrium statistical mechanics.
- Deal with time-dependent probabilities: stochastic equations.

Concluding Remarks

- What has all this to do with research at MSDL ?
- Look for **universal** features and quantify them: symmetry properties; conservation laws ?
- Equivalent concepts to: generalized coordinates, energy, energy minimization.
- Typical scales in a problem: length and time. Approximations based on this.
- Map one problem onto another: reduce it to something you know (SHO) .
- Apply the statistical mechanical approach to agents.